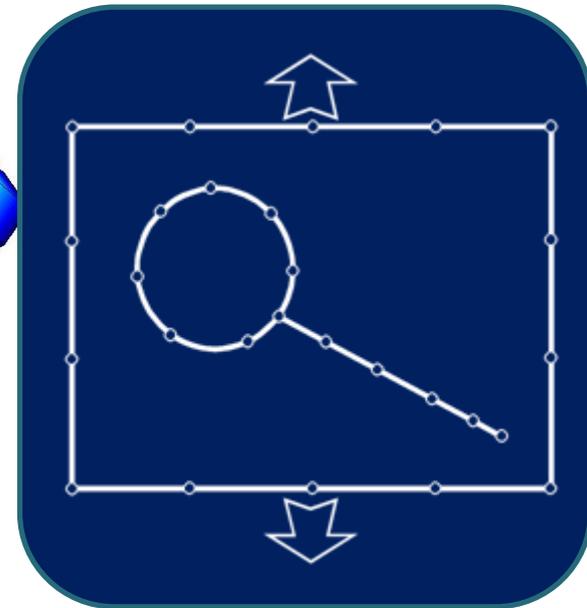
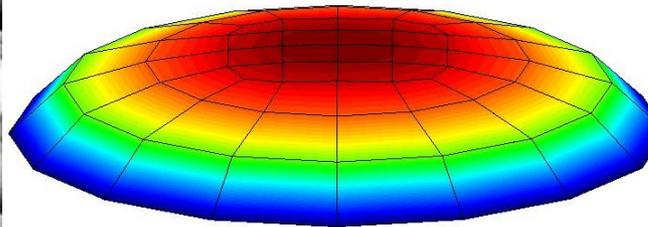


BEM for modeling fracture mechanics problems

Andrés Sáez

University of Seville (Spain)

Collaborators: J. Domínguez, Ch. Zhang, F. García-Sánchez, R. Rojas-Díaz, M. Wünsche



**NSF Workshop on the Emerging Applications and
Future Directions of the Boundary Element Method**

September 1-3, 2010

Cleveland, Ohio, USA (The University of Akron Campus)

Motivation:

Structural Integrity of Advanced Materials

Design of structural materials demands for high performance applications: we want lighter materials and new capacities

- ▣ Tailored materials with good strength/weight ratio → composites
- ▣ Multifield materials for control and smart structures applications → piezoelectricity...

Damage Tolerance Philosophy

- ▣ Existence of a flaw (crack...) in a mechanical component will no imply the end of its service life → need to evaluate its performance and reliability

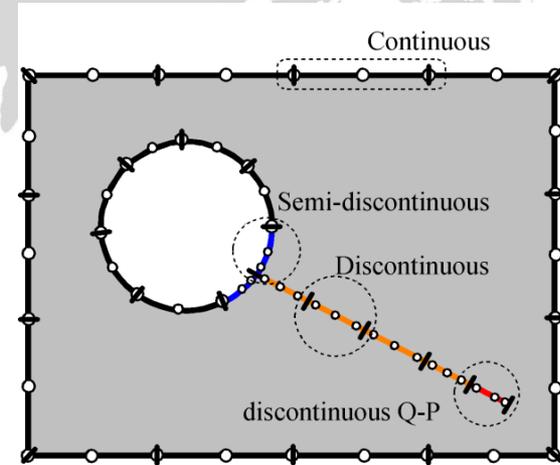
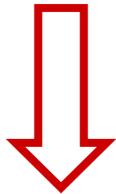


Motivation:

Structural Integrity of Advanced Materials

Tools to clearly characterize how a damaged structural element behaves are thus needed in order to predicts its:

- ▣ working conditions &
- ▣ remaining service life



Structural Integrity of Advanced Materials

The AIM is to extend to this new class of materials the damage estimation techniques previously developed for *classic* materials

WE'LL TACKLE THE PROBLEM FROM A NUMERICAL POINT OF VIEW (BEM)

Outline

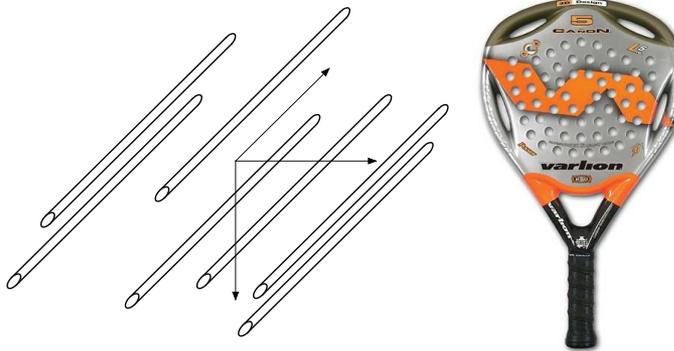
- Fracture mechanics of advanced materials:
 - Anisotropic materials (composites)
 - Multifield materials:
 - Piezoelectric
 - Magneto-electroelastic
- Why BEM for fracture?
- BEM strategies for fracture mechanics
- Dual or Hypersingular BEM
 - Statics
 - Dynamics
- Concluding remarks

Fracture mechanics of anisotropic materials

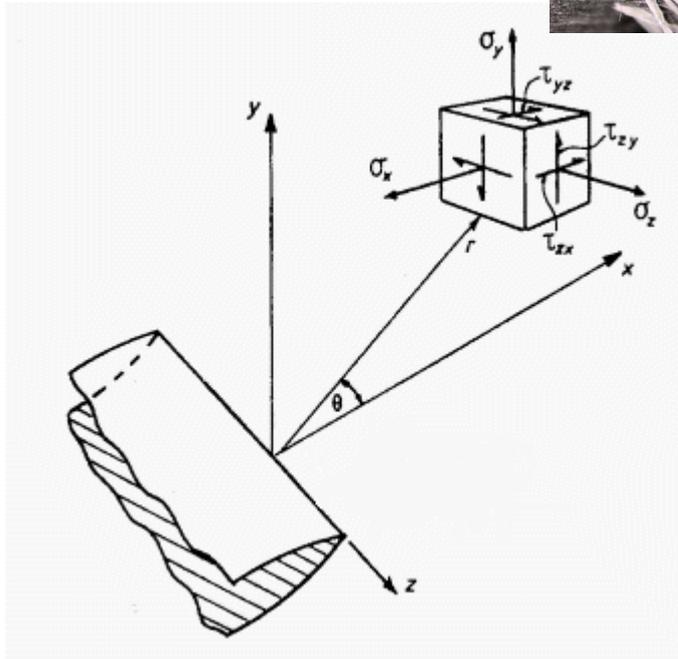
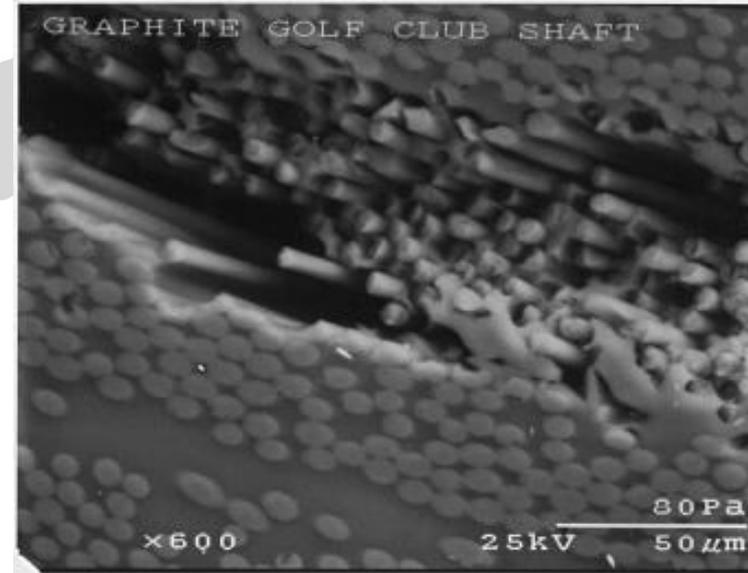
- Anisotropic materials in nature: Zinc, magnesium, wood, ice...
- Engineered materials: composites

2-D anisotropic behavior law

$$\begin{Bmatrix} \varepsilon_{11} \\ \varepsilon_{22} \\ 2\varepsilon_{12} \end{Bmatrix} = \begin{pmatrix} a_{11} & a_{12} & a_{16} \\ a_{12} & a_{22} & a_{26} \\ a_{16} & a_{26} & a_{66} \end{pmatrix} \begin{Bmatrix} \sigma_{11} \\ \sigma_{22} \\ \sigma_{12} \end{Bmatrix}$$

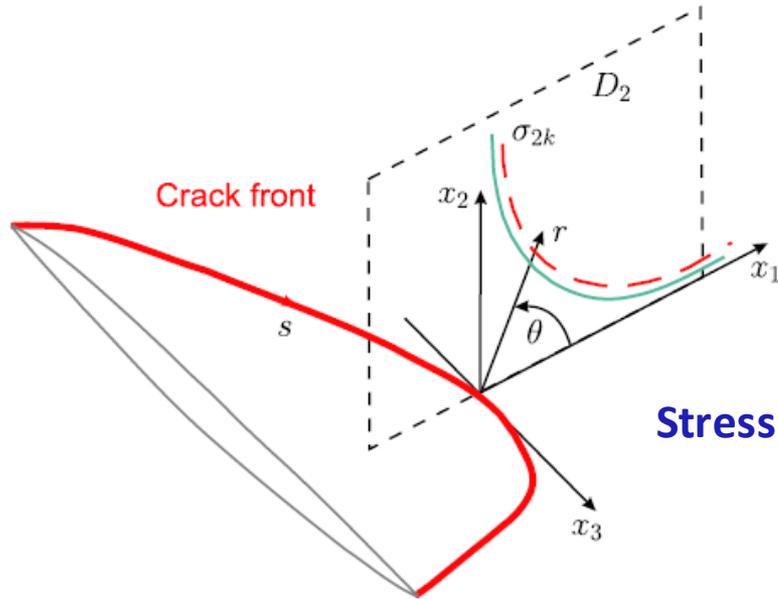


The existence of a crack will alter the stress fields and may lead to failure for loads much lower than the design load bearing capacity of the component



- MICRO-scale
- MESO-scale
- **MACRO-scale** ← This talk

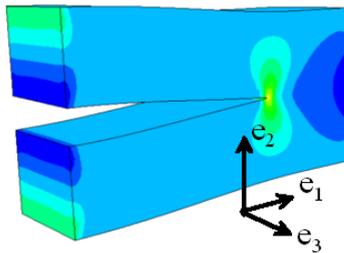
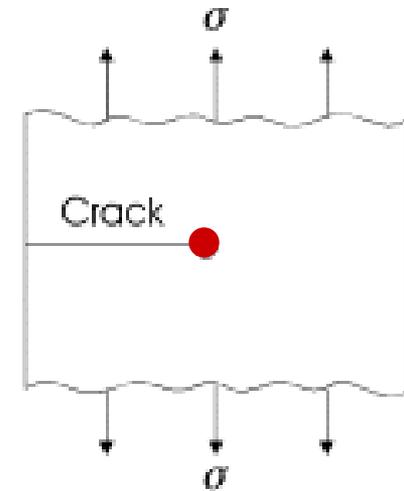
Displacement and stress fields around the crack tip⁽¹⁾:



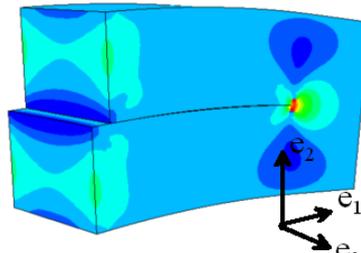
$$\sigma = K f(\theta, C_{ij}) / \sqrt{r}$$

$$u = K g(\theta, C_{ij}) \sqrt{r}$$

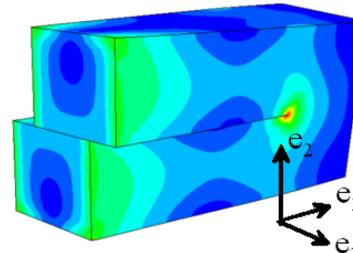
Stress Intensity Factors (SIF)



Mode I

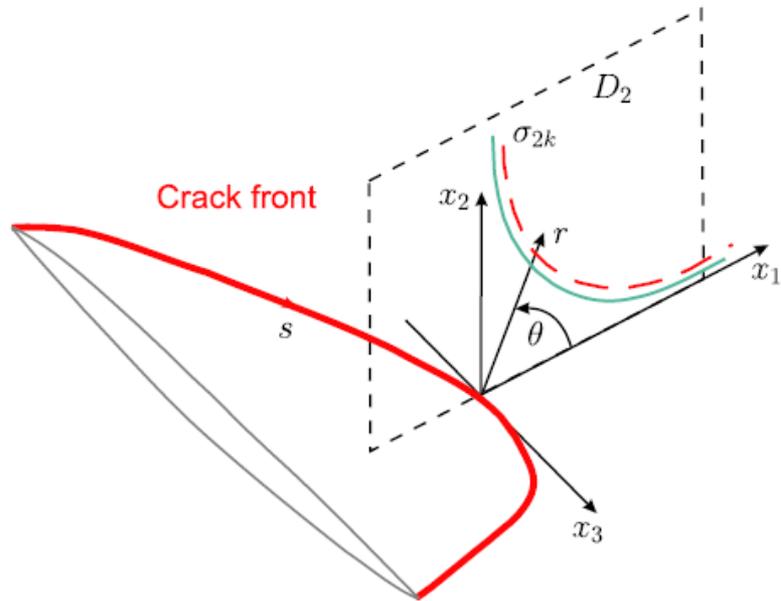


Mode II



Mode III

(1) Sih, Paris e Irwin (1965)



$$\sigma = K f(\theta, C_{ij}) / \sqrt{r}$$

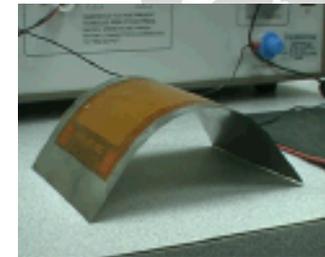
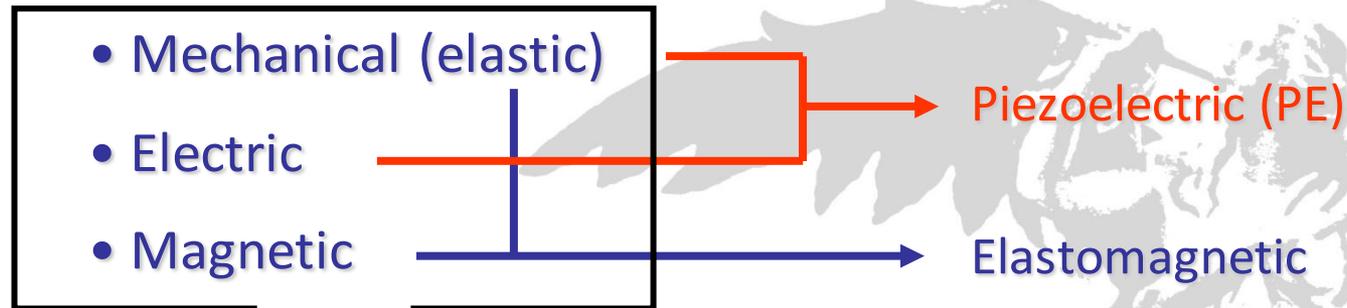
$$u = K g(\theta, C_{ij}) \sqrt{r}$$

When tackling the problem numerically, the selected method should fulfill some requirements:

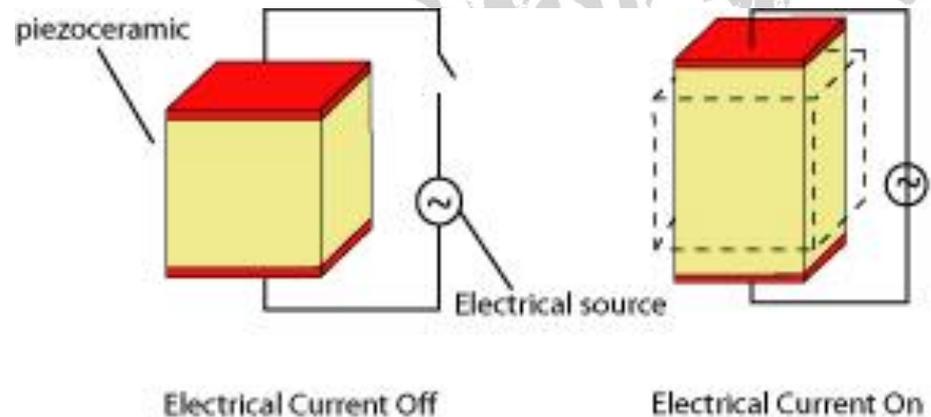
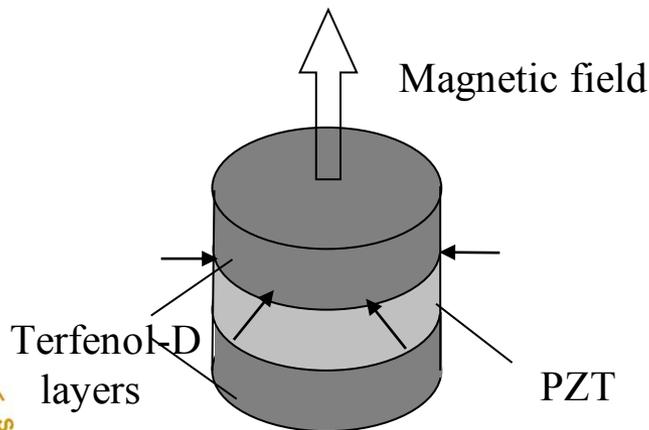
- ◆ Adequate representation of the mechanical fields around the crack
- ◆ SIF (or other fracture parameters) accurate computation
- ◆ Easy crack discretization

Fracture mechanics of multifield materials

Ability to **convert energy** among mechanical and non-mechanical fields:

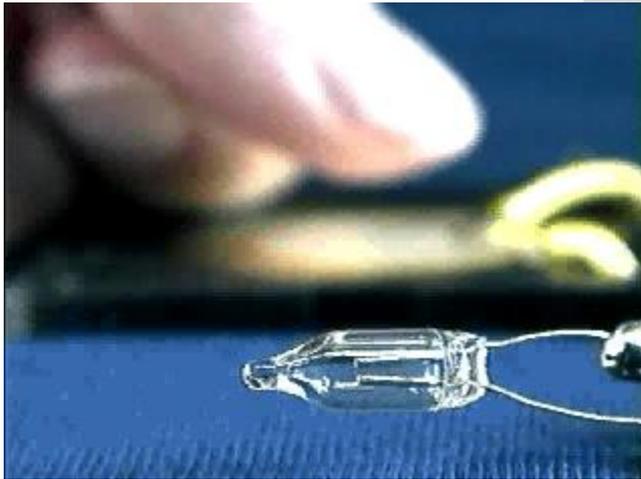


Magneto-electro-elastic (MEE)



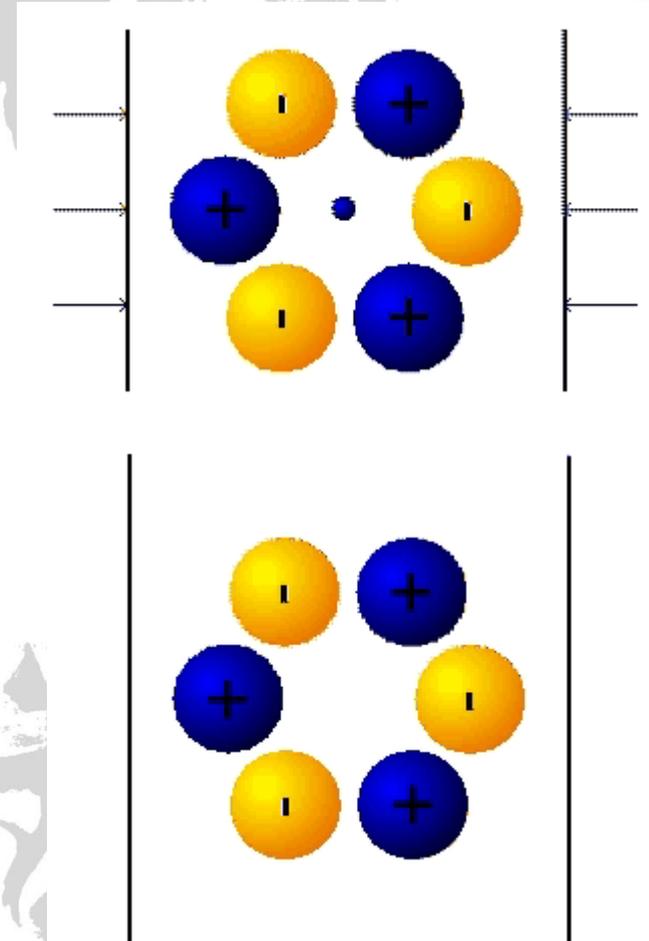
Piezoelectric Materiales

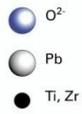
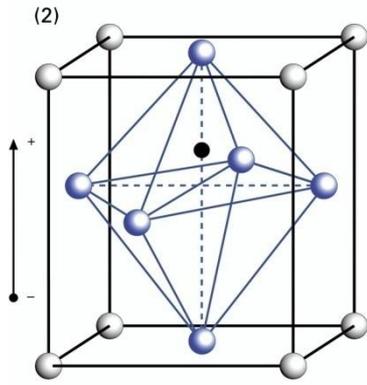
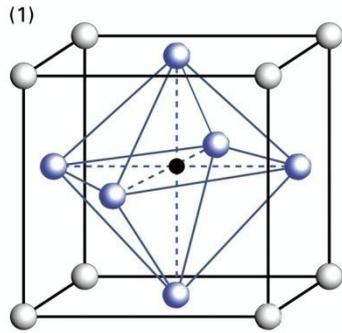
Direct piezoelectric effect ⁽¹⁾: a voltage is produced when the material is under tension or compression stress.



⁽¹⁾ Jacques y Pierre Curie, 1880

Inverse piezoelectric effect: when a potential difference is applied across the crystal it causes its deformation.



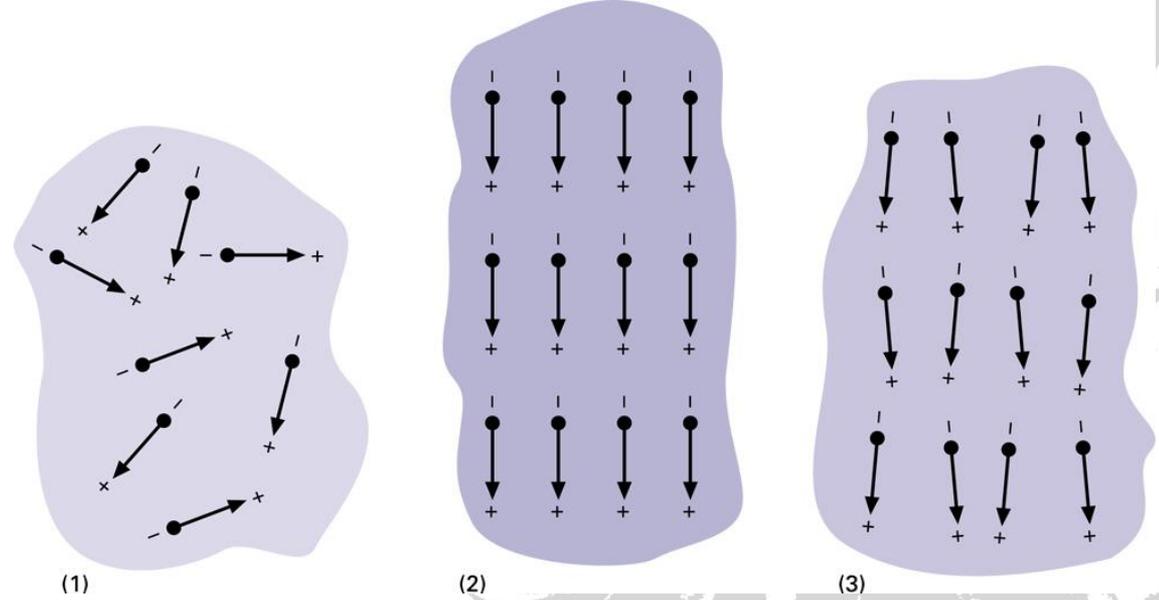


PZT unit cell:

1. PZT unit cell in the symmetric cubic state above the Curie temperature.
2. Tetragonally distorted unit cell below the Curie temperature.

PE materials

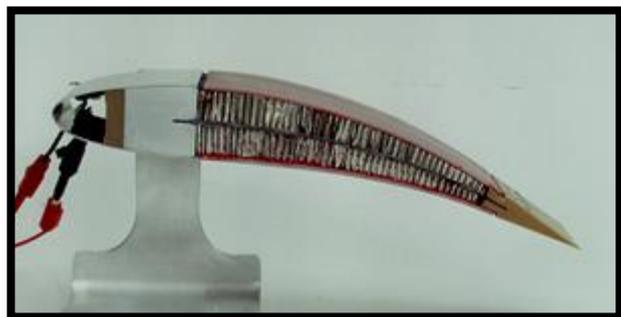
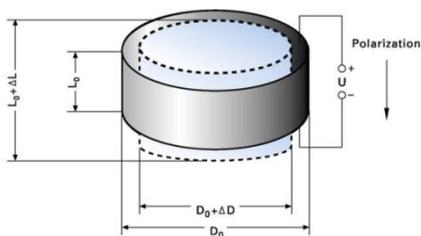
- ▣ Are always **anisotropic**
- ▣ Show larger PE coupling for artificial materials (e.g., **PZT ceramics**)



Electric dipoles in domains:

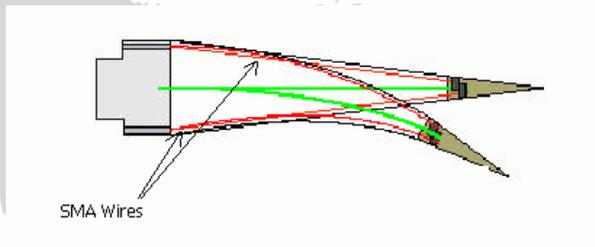
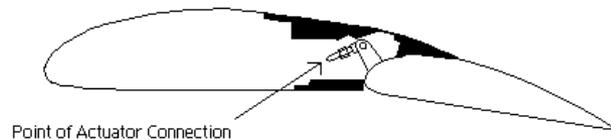
1. unpoled ferroelectric ceramic
2. during and
3. after poling (piezoelectric ceramic).

Applications for PE materials:

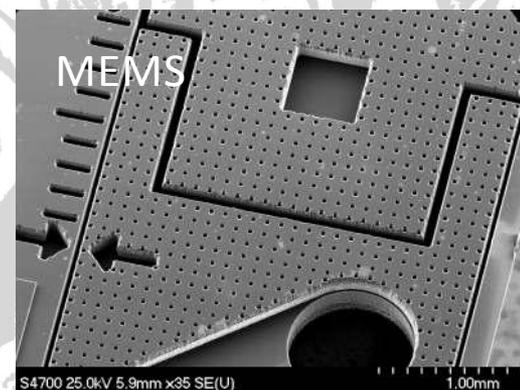
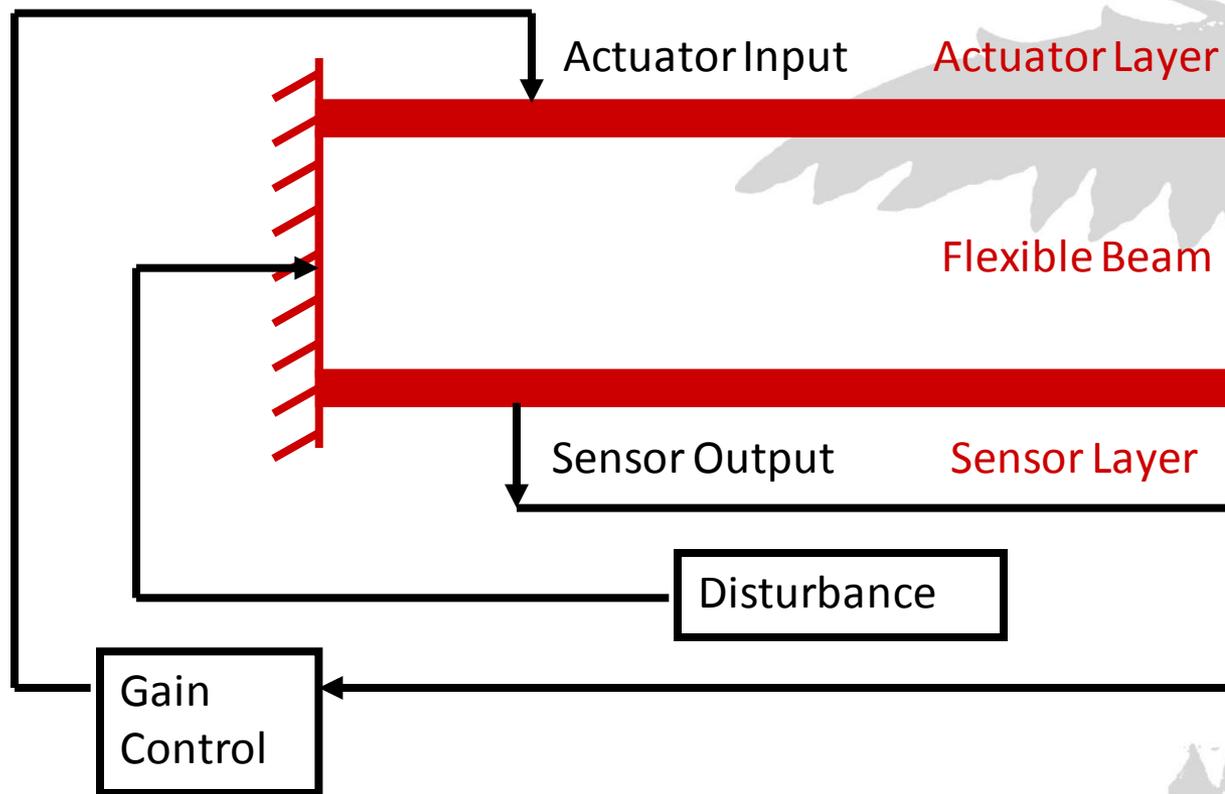


☐ sensors

☐ actuators



Applications for PE materials: Smart structures



How do we model PE materials?:

Extended notation for piezoelectricity (Barnett & Lothe, 1975)

Extended displacements vector

$$u_I = \begin{cases} u_i & I = 1,2,3 \\ \varphi & I = 4 \end{cases}$$

Electric potential

Extended stress tensor

$$\sigma_{iJ} = \begin{cases} \sigma_{ij} & I = 1,2,3 \\ D_i & I = 4 \end{cases}$$

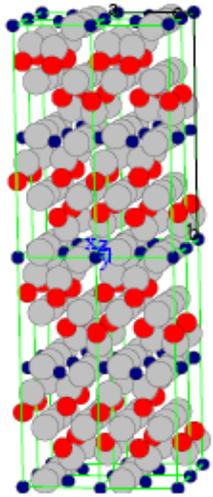
Electric displacement

PE behavior law

$$\begin{pmatrix} \sigma_{ij} \\ D_i \end{pmatrix} = \underbrace{\begin{pmatrix} C_{ijlm} & e_{kji} \\ e_{ijk} & -\kappa_{ij} \end{pmatrix}}_{C_{iJLm}} \begin{pmatrix} u_{l,m} + u_{m,l} \\ \varphi_{,j} \end{pmatrix} \quad \boxed{\sigma_{iJ,i} = 0}$$

elastic, piezoelectric & dielectric constants

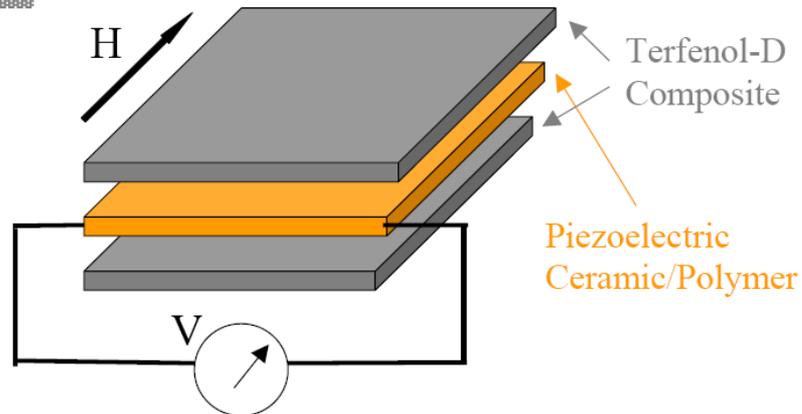
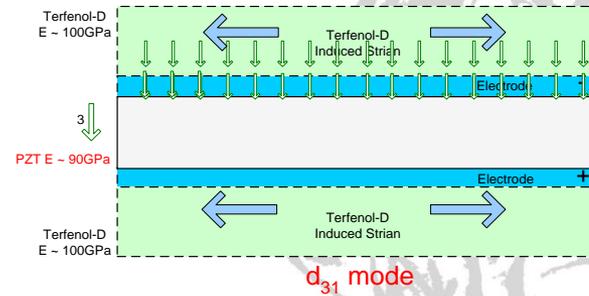
Magnetoelastic materials



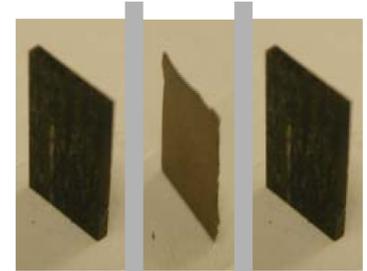
$YMnO_3$

$BiMnO_3$

Composite materials consisting of both piezoelectric and piezomagnetic phases may exhibit a magnetoelastic coupling effect that is not shown by any of the material phases alone. The magnetoelastic coupling of the resulting composite may be much larger than that of a single phase magnetoelastic material (Van Suchtelen, 1972; Nan, 1994; Benveniste, 1995).



Conductive Epoxy



Laminated [2-2] Magnetoelastic Composite

Terfenol-D Composite 5H PZT-5H Terfenol-D Composite



Modeling MEE materials

Magnetoelastic materials may be formulated in an elastic-like fashion by using the following generalized notation:

Extended displacement vector

$$u_I = \begin{cases} u_i & I = 1, 2, 3 \\ \phi & I = 4 \\ \varphi & I = 5 \end{cases}$$

Electric potential

Magnetic potential

Extended stress tensor

$$\sigma_{iJ} = \begin{cases} \sigma_{ij} & I = 1, 2, 3 \\ D_i & I = 4 \\ B_i & I = 5 \end{cases}$$

Electric displacement

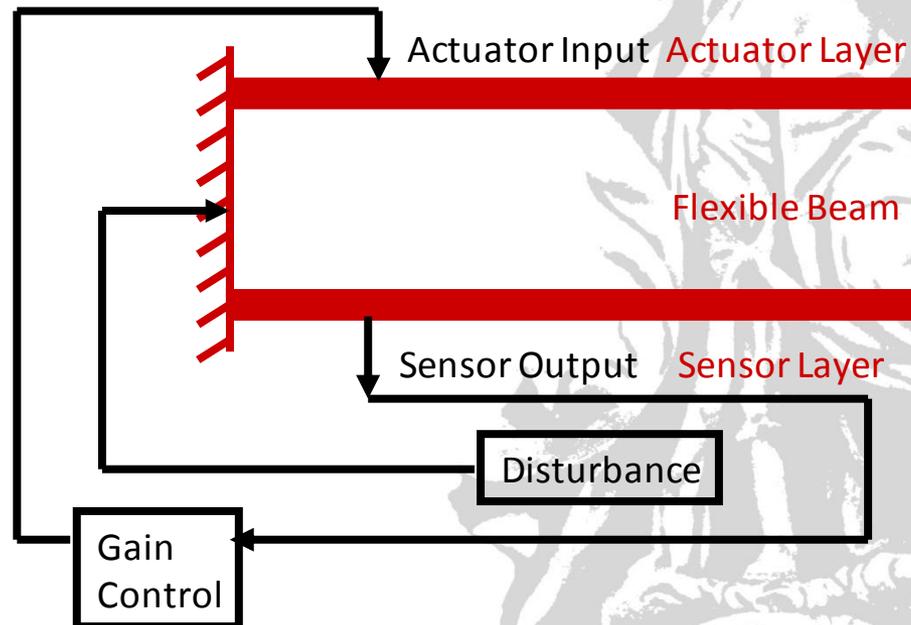
Magnetic induction

Constitutive equations:

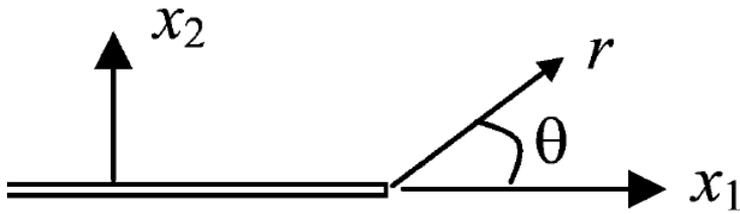
$$\sigma_{iJ} = C_{iJKl} u_{K,l}$$

Why study fracture in multifield materials?:

- structural integrity
- cracks alter the electric/magnetic reading: may lead to adopt wrong decisions (the structure is no longer smart)
- the electric and magnetic fields have influence on the crack growth process



Fracture mechanics of multifield materials



$$\sigma = K f(\theta, C_{ij}) / \sqrt{r}$$

$$u = K g(\theta, C_{ij}) \sqrt{r}$$

Generalized intensity factors:

- stress intensity factors (SIF): K_I and K_{II}
- electric displacement intensity factor (EDIF): K_D
- magnetic induction intensity factor (MIIF): K_B

Fracture of PE materials

Parton, V.Z., *Fracture mechanics of piezoelectric materials*,
Acta Astronautica, 3 (1976), pp. 671-683.

Pak, Y.E., *Crack extension force in a piezoelectric material*,
Journal of Applied Mechanics, Transactions ASME, 57
 (1990), pp. 647-653.

Fracture of MEE materials

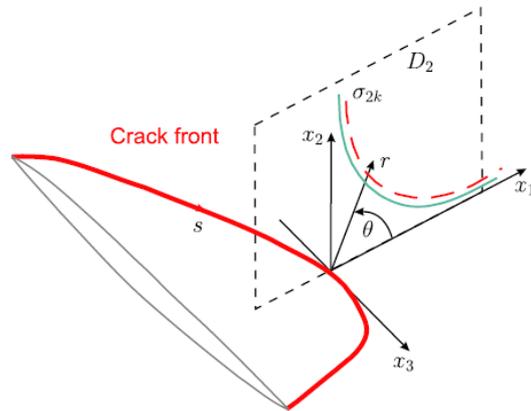
Song, Z.F., Sih, G.C., *Crack initiation behavior in
 magneto-electroelastic composite under in-plane deformation*,
Theor. and Appl. Fract. Mechanics, 39 (2003), pp. 189-207.

Wang, B.L., Mai, Y.-W., *Crack tip field in
 piezoelectric/piezomagnetic media*, *European Journal of
 Mechanics, A/Solids*, 22 (2003), pp. 591-602.



Why BEM for fracture?

Is BEM able to answer properly the questions we previously raised?



$$\sigma = K f(\theta, C_{ij}) / \sqrt{r}$$

$$u = K g(\theta, C_{ij}) \sqrt{r}$$

When tackling the problem numerically, the selected method should fulfill some requirements:

- ◆ Adequate representation of the field variables around the crack
- ◆ SIF (or other fracture parameters) accurate computation
- ◆ Easy crack discretization

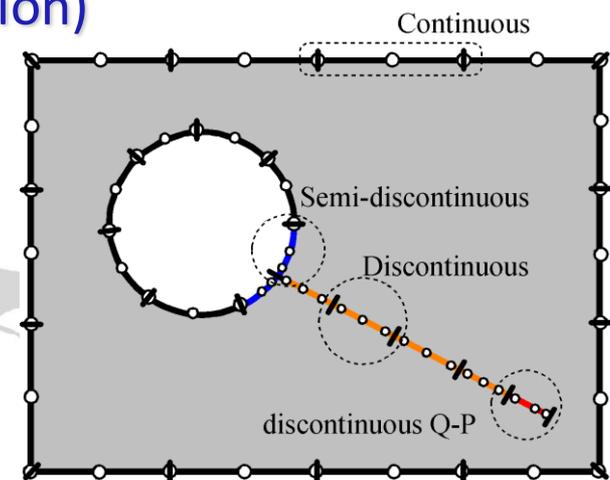
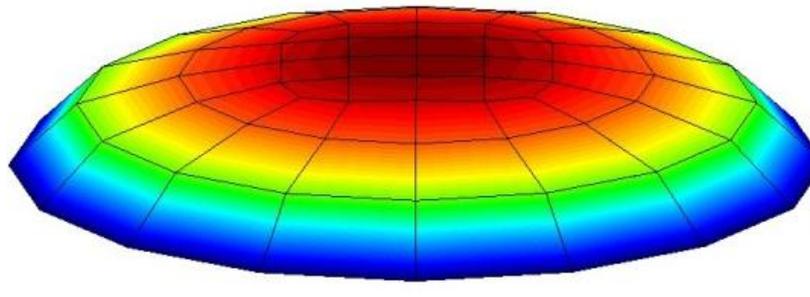
The answer is YEAH !!!!

Some BEM advantages

- The mesh is reduced in one dimension
- Automatic satisfaction of the radiation conditions at infinity
- Ability to capture high stress gradients
- Easy implementation of elements modeling crack-tip fields in fracture

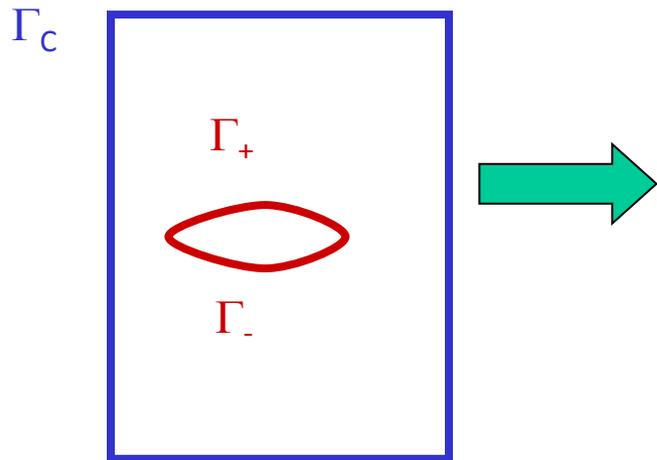
Some disadvantages

- Singular integrations
- (Need appropriate fundamental solution)



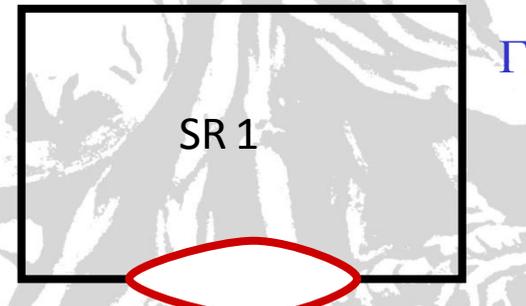
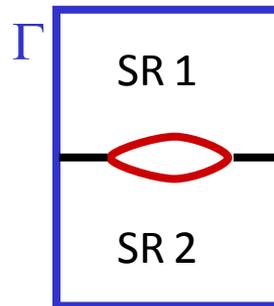
BEM strategies for fracture mechanics

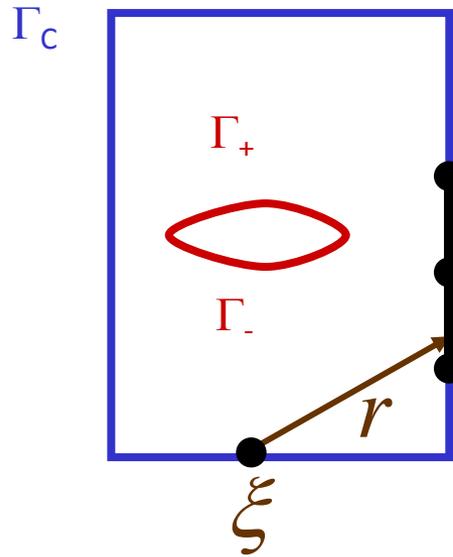
$$c_{IJ}u_J(\xi) + \int_{\Gamma} p_{IJ}^*(\xi, x)u_J(x)d\Gamma(x) = \int_{\Gamma} u_{IJ}^* p_J d\Gamma$$



0. Use *tailored* fundamental solutions

1. Subregion technique





2. Dual or hypersingular BEM

Displacement Boundary Integral Equation (BIE)

ξ Collocation point
 x Observation point

$$c_{IJ} u_J(\xi) + \int_{\Gamma} p_{IJ}^*(\xi, x) u_J(x) d\Gamma(x) = \int_{\Gamma} u_{IJ}^* p_J d\Gamma$$

$N_r C_{iJLr} \frac{\partial}{\partial \xi}$ Traction BIE

$$c_{IJ} p_J + N_r \int_{\Gamma} s_{rIJ}^* (u_J) d\Gamma = N_r \int_{\Gamma} d_{rIJ}^* p_J d\Gamma$$

$$u^* \propto \Theta(\ln r) \Rightarrow d^* \propto \Theta(1/r)$$

$$p^* \propto \Theta(1/r) \Rightarrow s^* \propto \Theta(1/r^2)$$

Extended C.O.D. $\begin{matrix} \Delta u \\ \Delta \phi \\ \Delta \varphi \end{matrix}$ $\Delta p = 0$
 $D = 0$
 $B = 0$

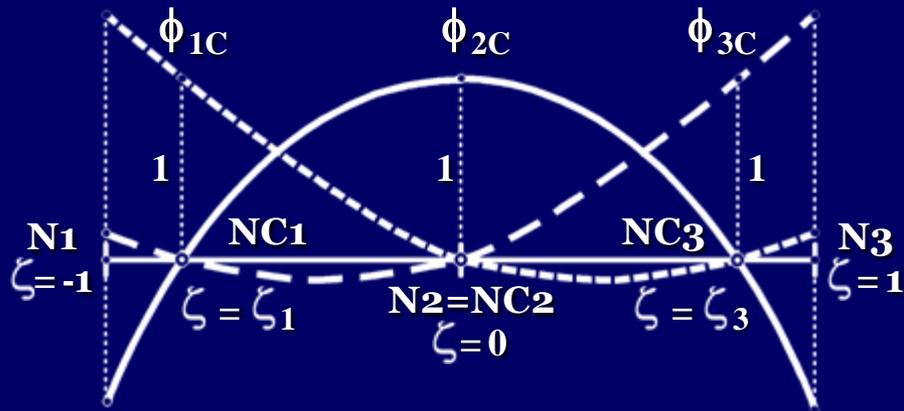


Dual or hypersingular BEM

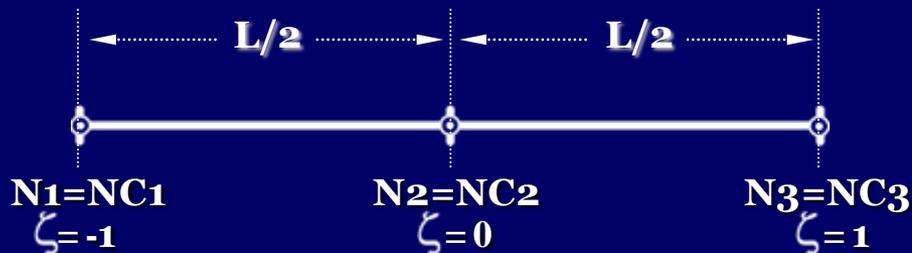
$$c_{IJ} u_J(\xi) + \int_{\Gamma} p_{IJ}^*(\xi, x) u_J(x) d\Gamma(x) = \int_{\Gamma} u_{IJ}^* p_J d\Gamma$$

$$c_{IJ} p_J + N_r \int_{\Gamma} s_{rIJ}^* u_J d\Gamma = N_r \int_{\Gamma} d_{rIJ}^* p_J d\Gamma$$

- ❖ Meshing strategy?
- ❖ Availability of reasonable fundamental solutions for statics?
- ❖ Singular and hypersingular integrations?
- ❖ Computation of fracture parameters?
- ❖ What about dynamics?



continuous DBIE (Γ_c)



Traction BIE (TBIE) \rightarrow displacement field C^1

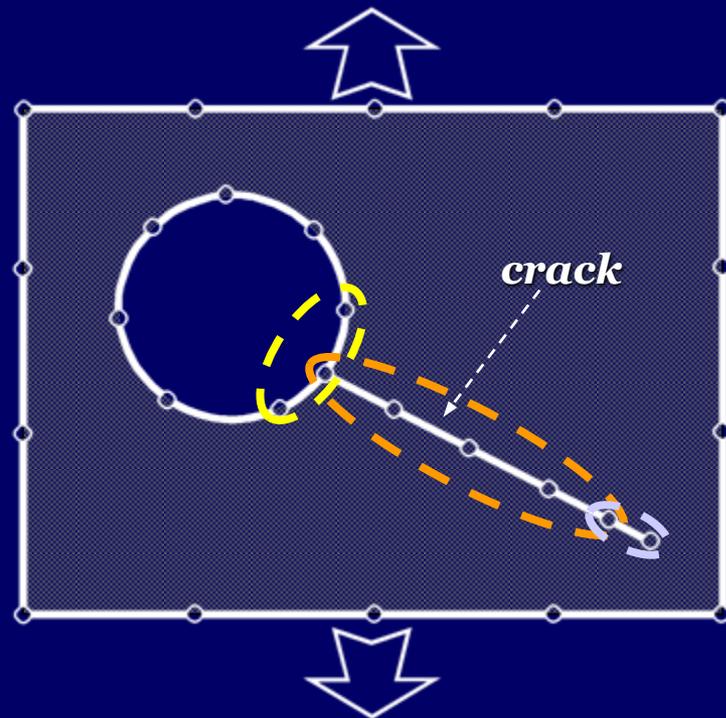
discontinuous TBIE (Γ_+)



semidiscontinuous EID (Γ_c)



Meshing strategy



discontinuous quarter-point (Γ_+)



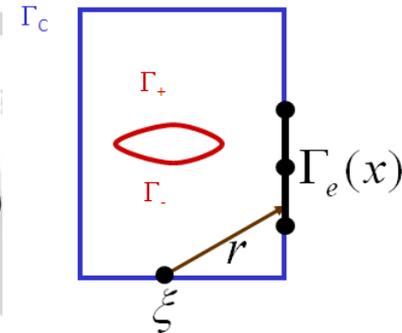
Static fundamental solution

$$u_{IJ}^{*S} = \text{Re} \left[\sum_{R=1}^5 A_{JR} H_{RI} \ln(z_R^x - z_R^\xi) \right] \rightarrow d_{IJ}^{*S} = d_{rIJ}^{*S} N_r \propto \text{Re} \left[\frac{\mu_R N_1 - N_2}{(z_R^x - z_R^\xi)} \right]$$

$$p_{IJ}^{*S} = \text{Re} \left[\sum_{R=1}^5 L_{JR} H_{RI} \frac{\mu_R n_1 - n_2}{(z_R^x - z_R^\xi)} \right] \rightarrow s_{rIJ}^{*S} \propto \text{Re} \left[\frac{\mu_R n_1 - n_2}{(z_R^x - z_R^\xi)^2} \right]$$

Observation point ξ : $z_R^\xi = \xi_1 + \mu_R \xi_2 \iff \xi = (\xi_1, \xi_2)$

Collocation point \mathbf{x} : $z_R^x = x_1 + \mu_R x_2 \iff \mathbf{x} = (x_1, x_2)$

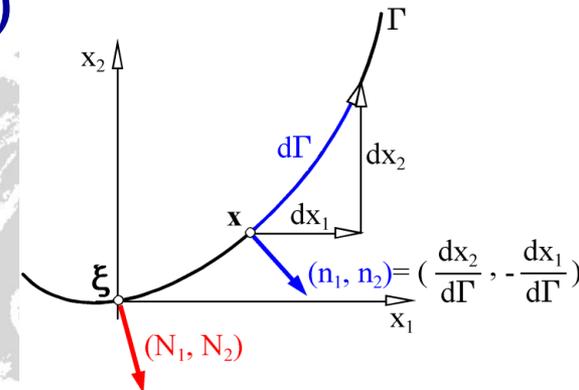


$$c_{IJ} u_J(\xi) + \int_{\Gamma} p_{IJ}^* u_J d\Gamma(\mathbf{x}) = \int_{\Gamma} u_{IJ}^* p_J d\Gamma(\mathbf{x})$$

Anisotropic (elastic): Eshelby et al. (1953)

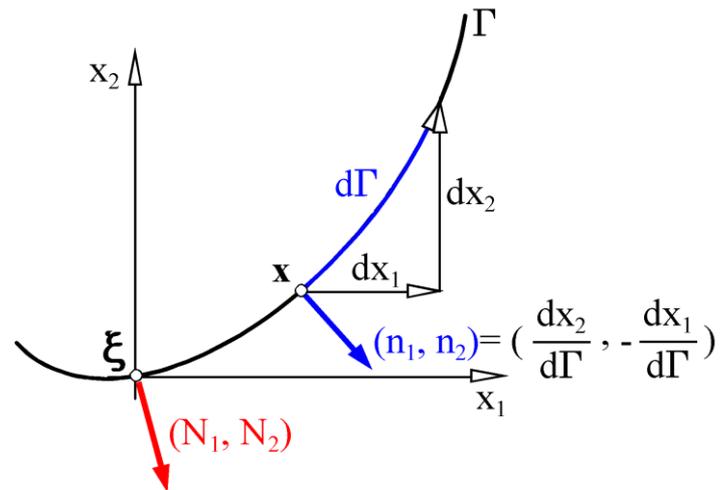
Piezoelectric: Barnett & Lothe (1975)

MEE: Liu et al. (2001)



Evaluation of Strongly Singular and Hypersingular Integrals

Regularization technique follows the works by García-Sánchez et al. (EABE2004, C&S2005, TAFM2008). This approach is valid for any of the 2-D fundamental solutions derived by Stroh's formalism and has no restrictions on the type/shape of the boundary elements:



$$p_{IJ}^*(\mathbf{x}, \boldsymbol{\xi}) = \frac{1}{\pi} \Re \left[L_{JM} Q_{MI} \frac{\mu_M n_1 - n_2}{z_M^x - z_M^\xi} \right]$$

CHANGE OF VARIABLES

$$\chi_R = z_R^x - z_R^\xi = (x_1 - \xi_1) + \mu_R (x_2 - \xi_2)$$

THE JACOBIAN IS BUILT-IN THE FUNDAMENTAL SOLUTION

$$\frac{d\chi_R}{d\Gamma} = \frac{d\chi_R}{dx_1} \frac{dx_1}{d\Gamma} + \frac{d\chi_R}{dx_2} \frac{dx_2}{d\Gamma} = -n_2 + \mu_R n_1$$

For instance, for the hypersingular integrations:

$$S_{rIJ}^* \approx O \left[\frac{1}{(z^x - z^\xi)^2} \right]$$

$$\frac{d\chi_R}{d\Gamma} = \frac{d\chi_R}{dx_1} \frac{dx_1}{d\Gamma} + \frac{d\chi_R}{dx_2} \frac{dx_2}{d\Gamma} = -n_2 + \mu_R n_1$$

$$I''_R = \int_{\Gamma_e} (\mu_R n_1 - n_2) \frac{1}{\chi_R^2} \phi_q d\Gamma$$

$$\begin{aligned} \phi_q(\chi_R) &\approx \phi_q(\chi_R = 0) + \left. \frac{d\phi_q}{d\chi_R} \right|_{\chi_R=0} \cdot \chi_R \\ &\approx \phi_{q0} + \phi'_{q0} \cdot \chi_R \end{aligned}$$

$$I''_R = \underbrace{\int_{\Gamma_e} \frac{1}{\chi_R^2} (\phi_q - \phi_{q0} - \phi'_{q0} \cdot \chi_R) d\chi_R}_{\text{Regular}} + \underbrace{\phi_{q0} \int_{\Gamma_e} \frac{1}{\chi_R^2} d\chi_R}_{\text{Analytic}} + \underbrace{\phi'_{q0} \int_{\Gamma_e} \frac{1}{\chi_R} d\chi_R}_{\text{Analytic}}$$

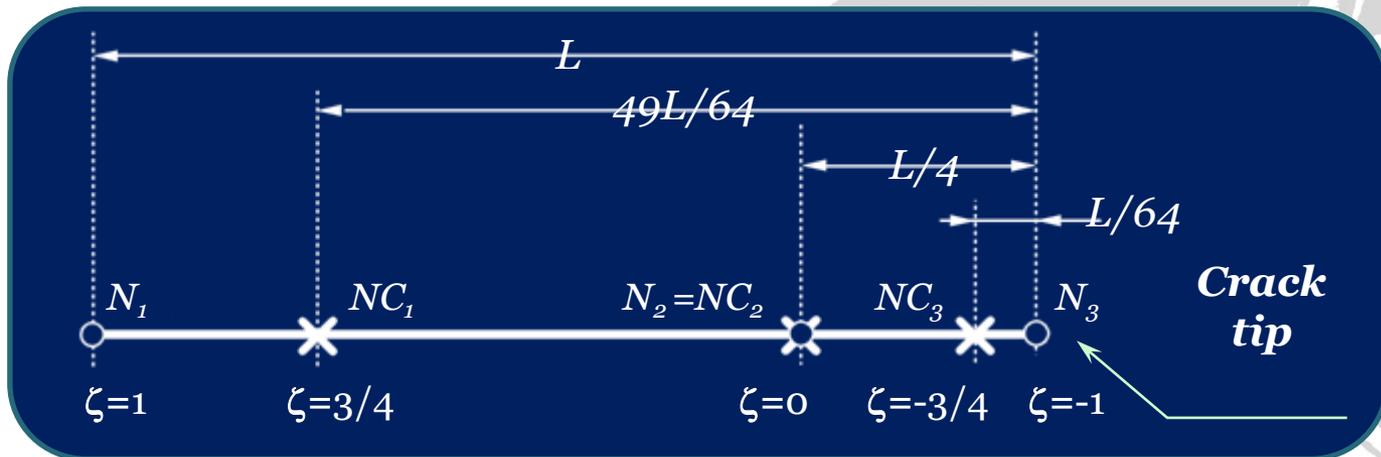
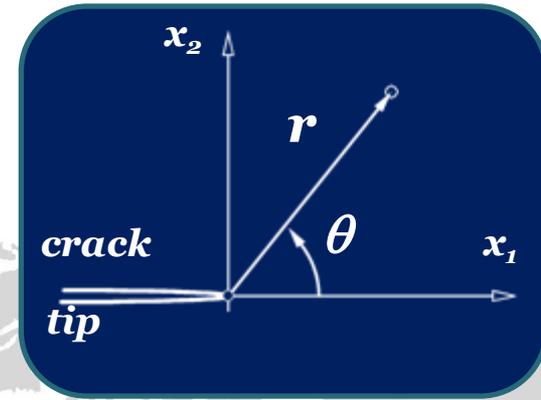
Regular

Analytic

Analytic

Evaluation of fracture parameters by extrapolation:

$$\Delta u \rightarrow K_* f_u(\theta, \mu_R) \sqrt{r}$$



$$\zeta = 2\sqrt{\frac{r}{L}} - 1$$

at NC_3

$$\mathbf{K} = \begin{pmatrix} K_{II} \\ K_I \\ K_D \\ K_B \end{pmatrix} = \sqrt{\frac{\pi}{8r}} \mathbf{H}^{-1} \begin{pmatrix} \Delta u_1 \\ \Delta u_2 \\ \Delta \phi \\ \Delta \varphi \end{pmatrix}$$

$$G = \frac{1}{2} \mathbf{K}^T \mathbf{H} \mathbf{K}$$



Some results: BEM vs FEM & Experimental Results. Composite material.



Available online at www.sciencedirect.com



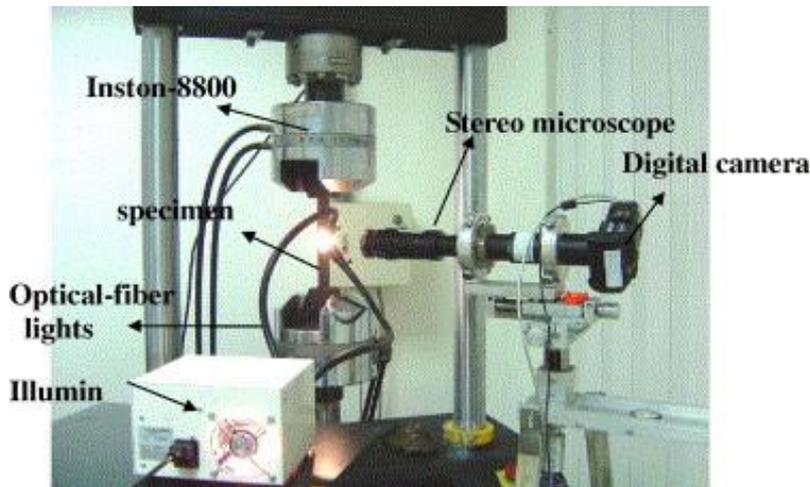
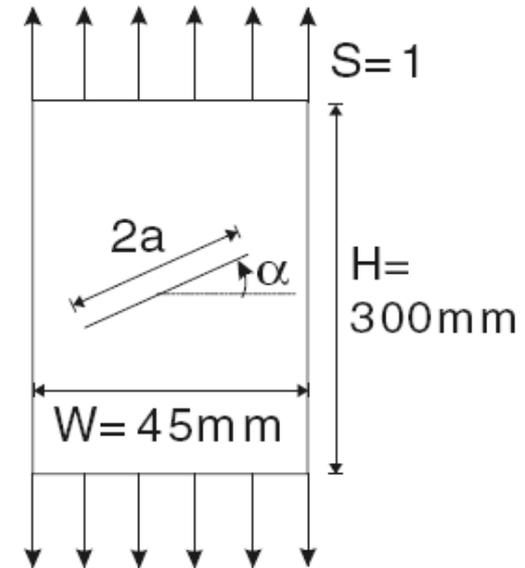
Composite Structures 81 (2007) 614-621

COMPOSITE
STRUCTURES

www.elsevier.com/locate/compstruct

Determining stress intensity factors of composites using crack opening displacement

S.H. Ju *, S.H. Liu

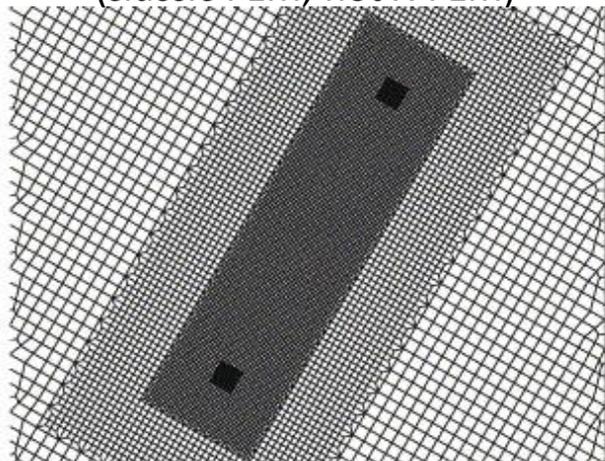


tip. All numerical analysis assume linear-elastic and plane-stress conditions. The applied stress, S , in each case is unity and the material properties are $E_{11} = 101$ GPa, $E_{22} = 25$ GPa, $\nu_{21} = 0.082$ and $G_{12} = 4.2$ Gpa, where fiber axis 1 is parallel to the specimen longitudinal direction and axis 2 is parallel to the transverse direction. As shown in

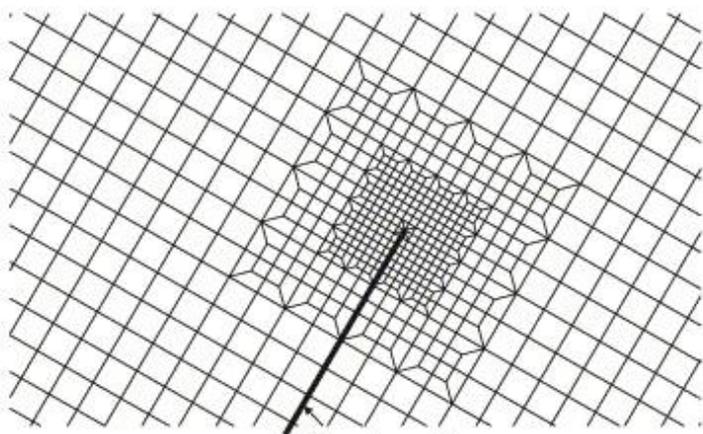
60° with respect to the horizontal. The material was a $[0_3/90/0_3/90/0_3/90/0_3]$ 12 K-carbon-fiber/epoxy composite.



FEM Mesh
(classic FEM, not X-FEM)

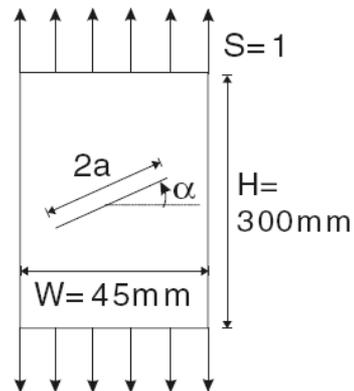


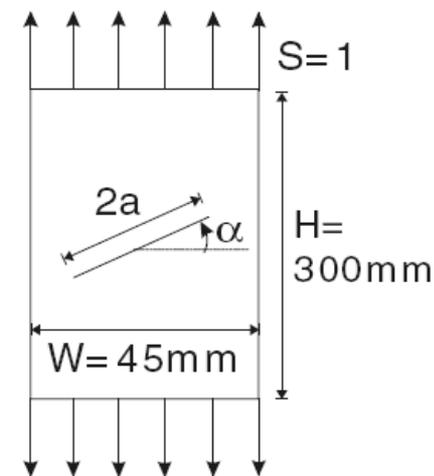
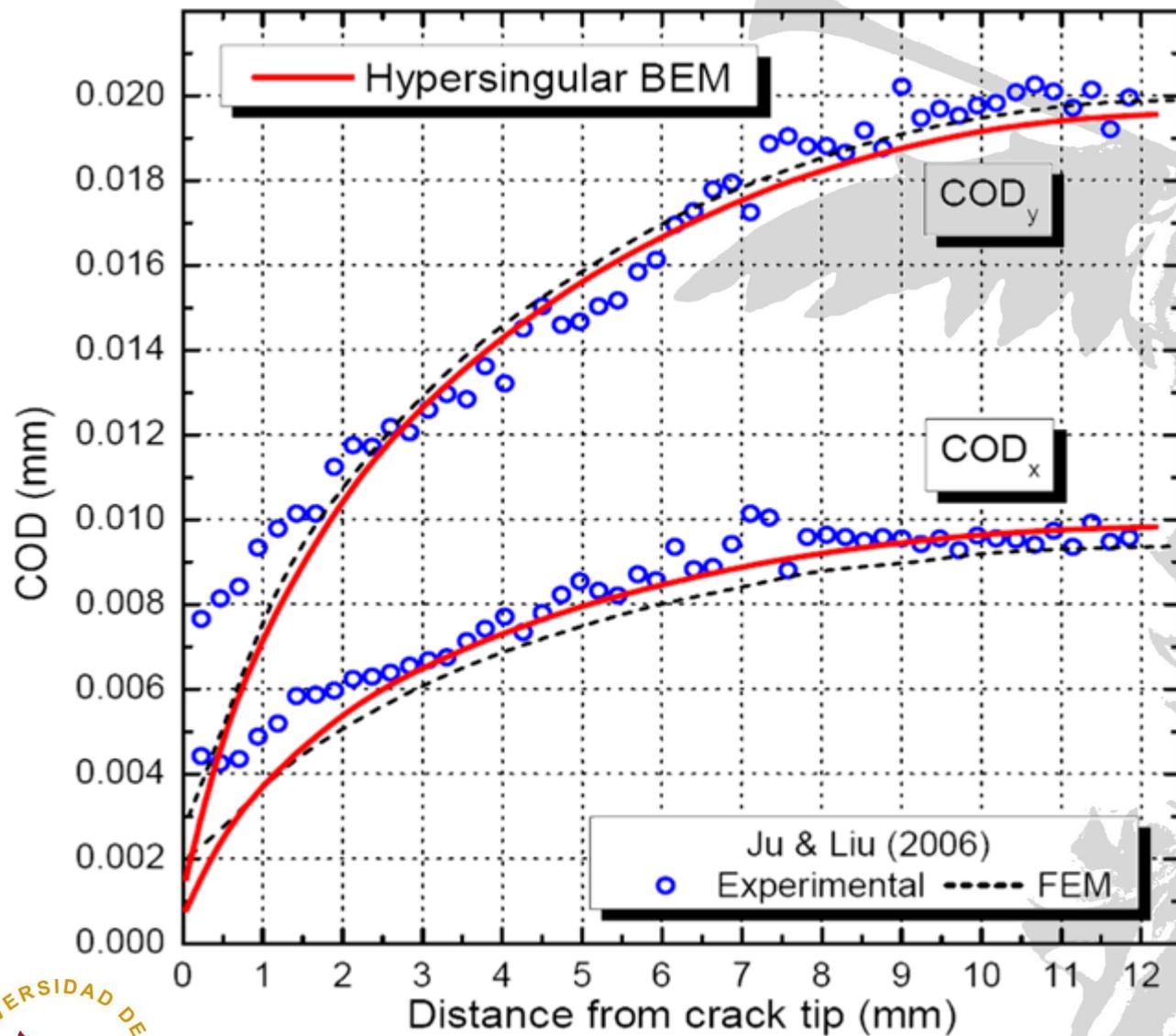
(a) Finite element mesh of near the crack of $\alpha=60^\circ$ and $2a/W=0.556$



(b) Finite element mesh very near the upper right crack tip

BEM Mesh





Some results: BEM vs FEM & Analytic. Piezoelectric material.

Arch Appl Mech (2006) 76: 725-745
 DOI 10.1007/s00419-006-0059-z

ORIGINAL

MEF

M. Kuna

Finite element analyses of cracks in piezoelectric structures: a survey

barium titanate (BaTiO_3)

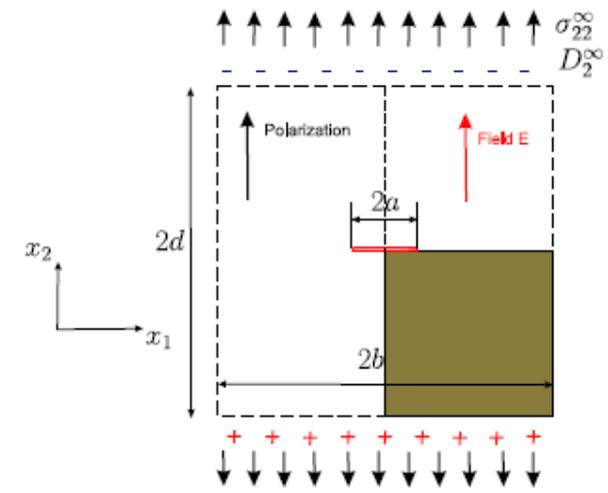


Fig. 9 Griffith crack under remote electromechanical loading

	Impermeable $\kappa_C \rightarrow 0$
K_I [$\text{MPa}\sqrt{\text{m}}$] exact	3.545
K_I (CTE) error	1.98%
K_I (MCCI) error	0.68%
K_{IV} [$\text{Cm}^{-3/2}$] exact	3.545E-03
K_{IV} (CTE) error	0.85%
K_{IV} (MCCI) error	0.37%

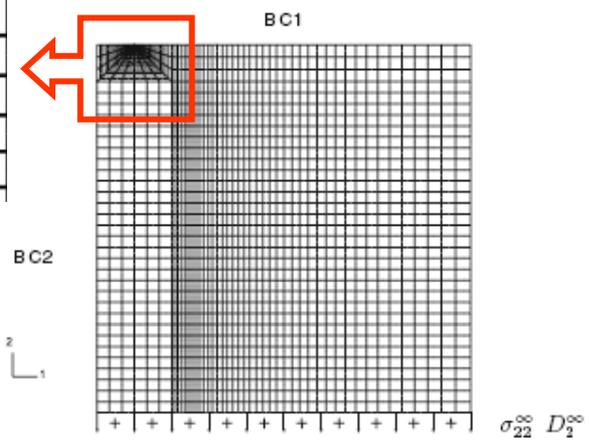
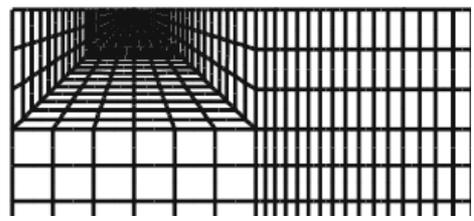
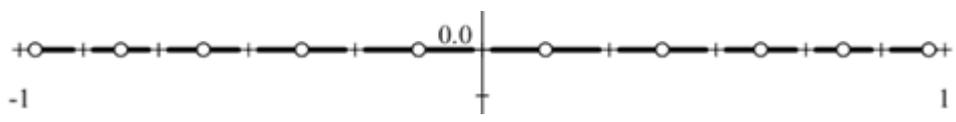


Fig. 10 FEM model and boundary conditions for Griffith crack

BEM

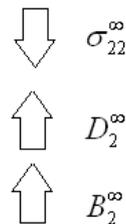
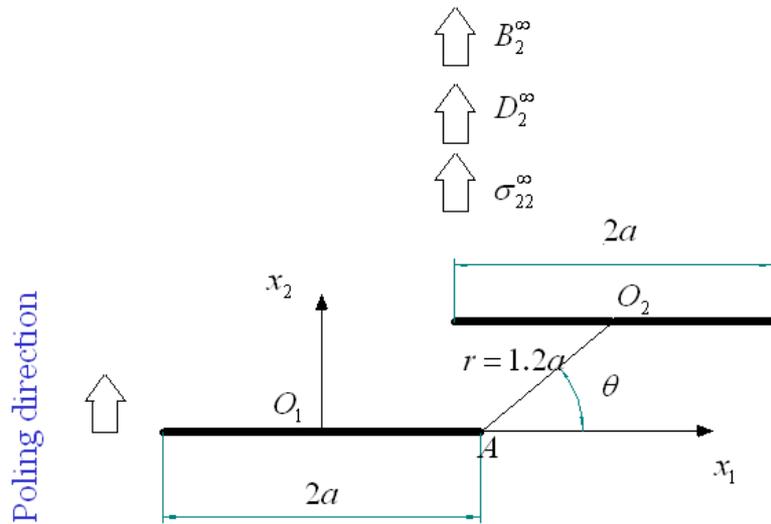


error < 0.1 %

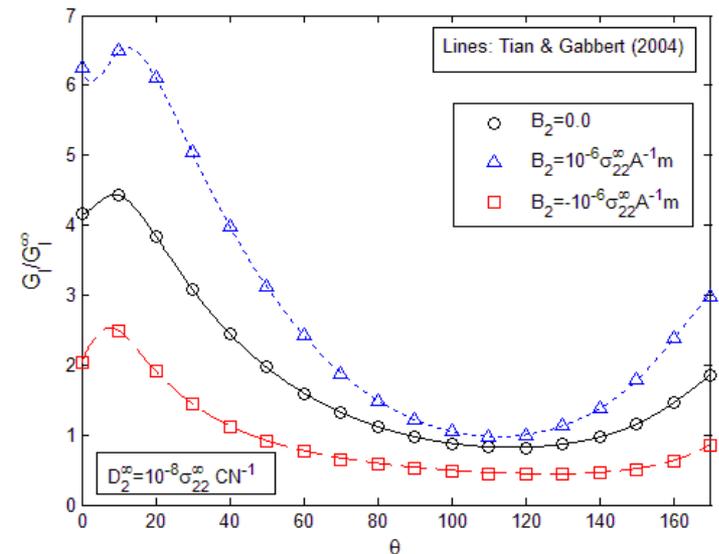
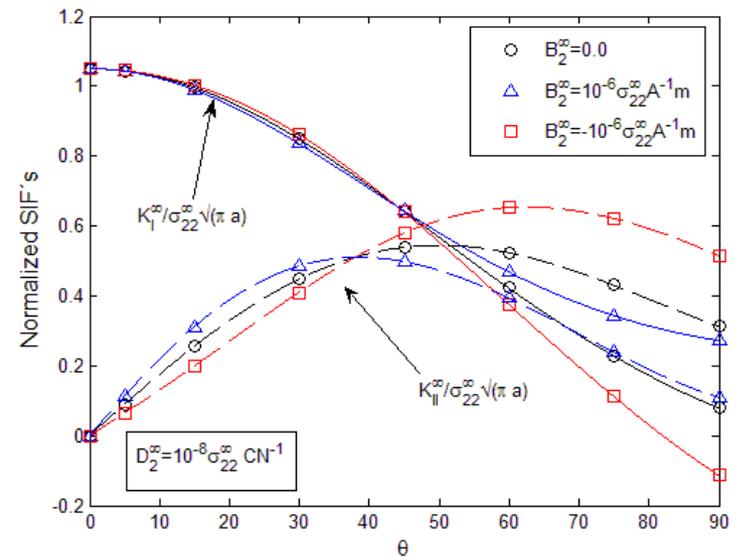


A. Sáez
 UNIVERSIDAD DE SEVILLA

Some results: BEM vs Analytic. MEE material.



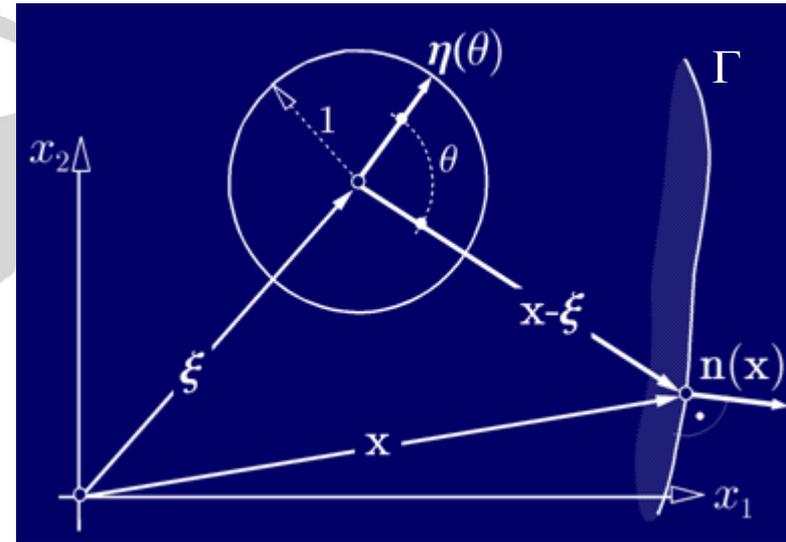
BaTiO₃
50%
CoFe₂O₄



What about dynamics?

For instance, in the frequency domain, the fundamental solution is obtained as:

$$u_{IJ}^*(\mathbf{x}, \boldsymbol{\xi}, \omega) = \frac{1}{8\pi^2} \int_{|\boldsymbol{\eta}|=1} \frac{\gamma_{IJ}^m}{\rho c_m^2} \underbrace{\Phi(k_m(\omega) |\boldsymbol{\eta} \cdot (\mathbf{x} - \boldsymbol{\xi})|)}_{\text{regular}} dS(\boldsymbol{\eta})$$



$$\underbrace{\frac{1}{8\pi^2} \int_{|\boldsymbol{\eta}|=1} \frac{\gamma_{IJ}^m}{\rho c_m^2 E_{qq}^m} [\Phi(k_m(\omega) |\boldsymbol{\eta} \cdot (\mathbf{x} - \boldsymbol{\xi})|) + 2 \ln |\boldsymbol{\eta} \cdot (\mathbf{x} - \boldsymbol{\xi})|]}_{\text{regular}} dS(\boldsymbol{\eta}) - \frac{1}{8\pi^2} \int_{|\boldsymbol{\eta}|=1} \frac{\gamma_{IJ}^m}{\rho c_m^2 E_{qq}^m} 2 \ln |\boldsymbol{\eta} \cdot (\mathbf{x} - \boldsymbol{\xi})| dS(\boldsymbol{\eta})}_{\text{singular (statics)}}$$

regular: $u_{IJ}^{*R}(\mathbf{x}, \boldsymbol{\xi}, \omega)$

singular:
(statics) $u_{IJ}^{*S}(\mathbf{x}, \boldsymbol{\xi})$

Anisotropic: Wang & Achenbach (1995)

Piezoelectric: Denda et al. (2004)

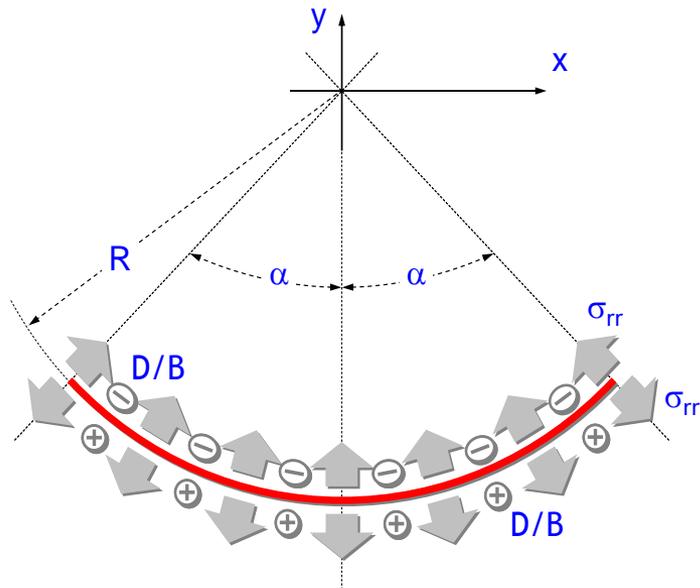
MEE: Rojas-Díaz et al. (2008)

$$u_{IJ}^*(\mathbf{x}, \boldsymbol{\xi}, \omega) = u_{IJ}^{*S}(\mathbf{x}, \boldsymbol{\xi}) + u_{IJ}^{*R}(\mathbf{x}, \boldsymbol{\xi}, \omega)$$



Some dynamic results: Combined magneto-electro-mechanical impacts.

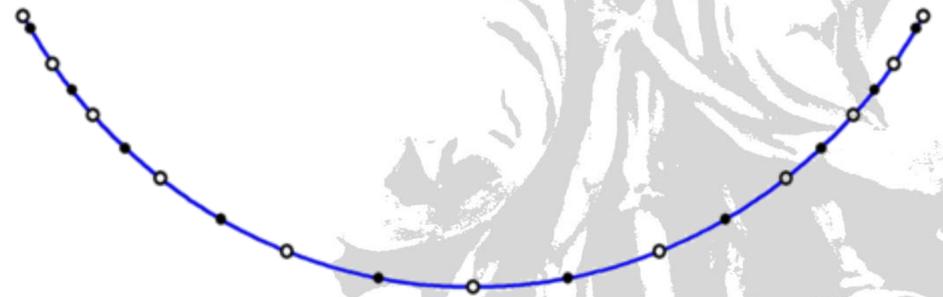
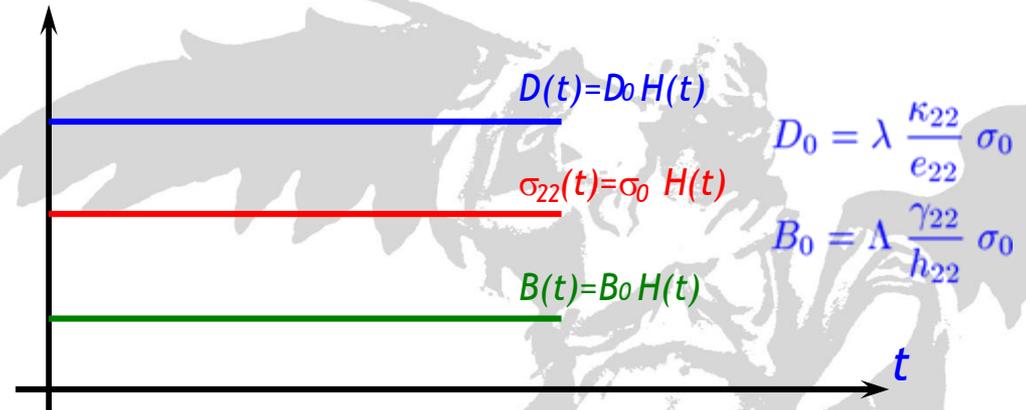
Curved crack in infinite MEE domain



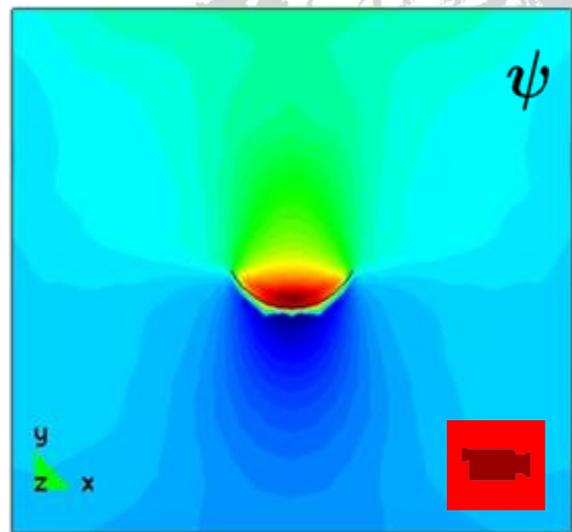
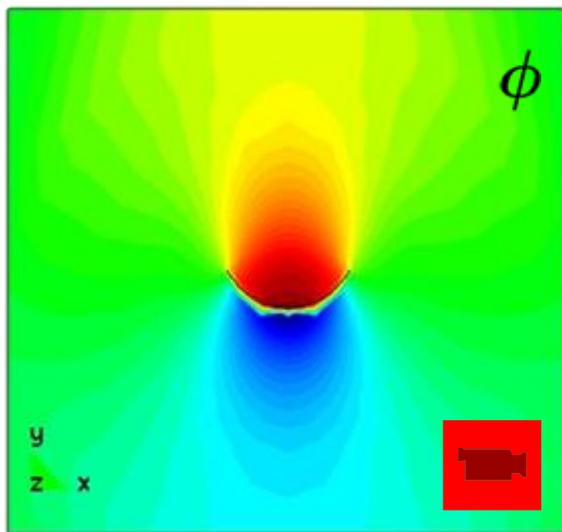
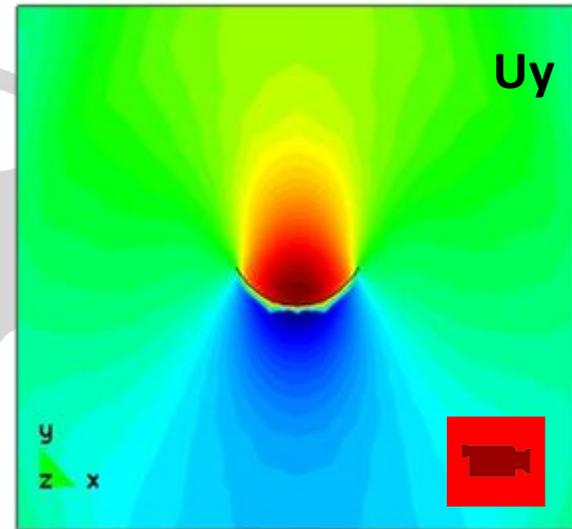
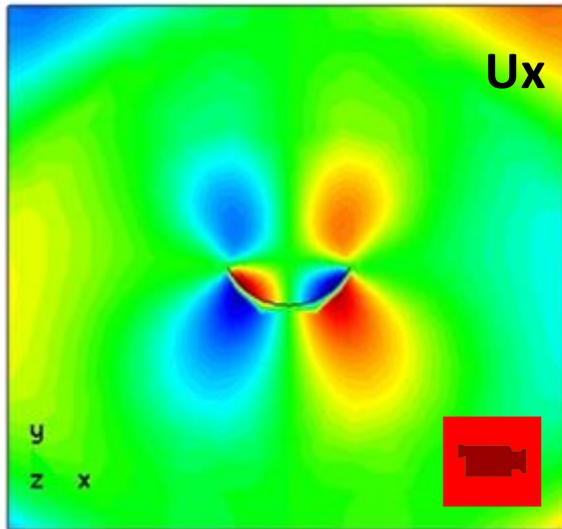
BaTiO₃

50%

CoFe₂O₄

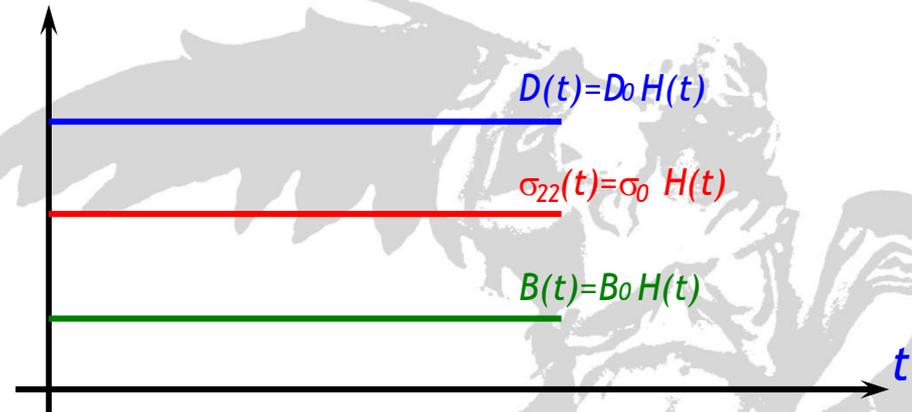


10 discontinuous quadratic elements for the crack (8 + 2 QP)

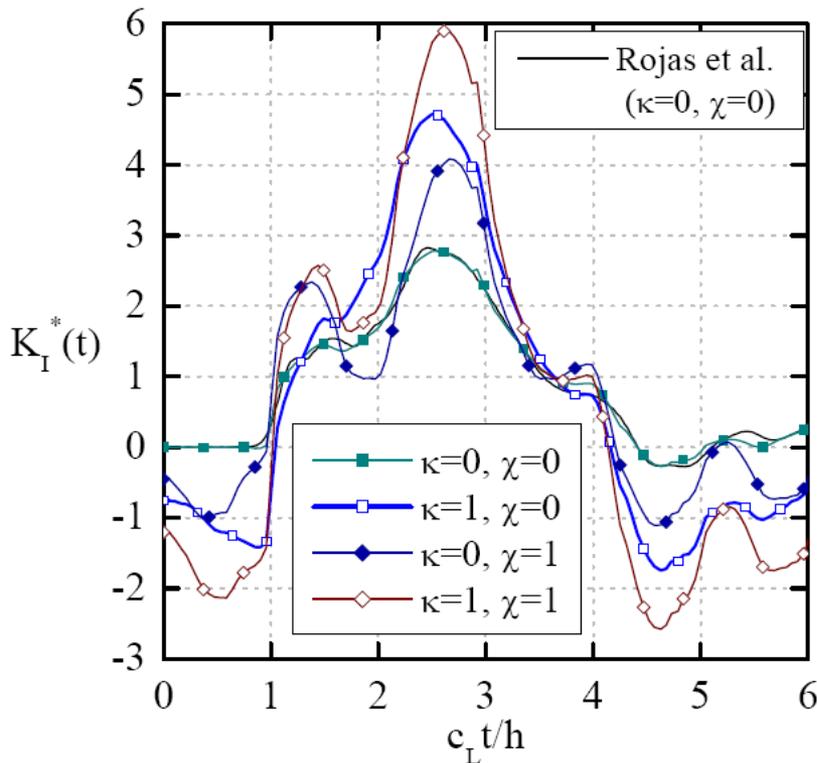
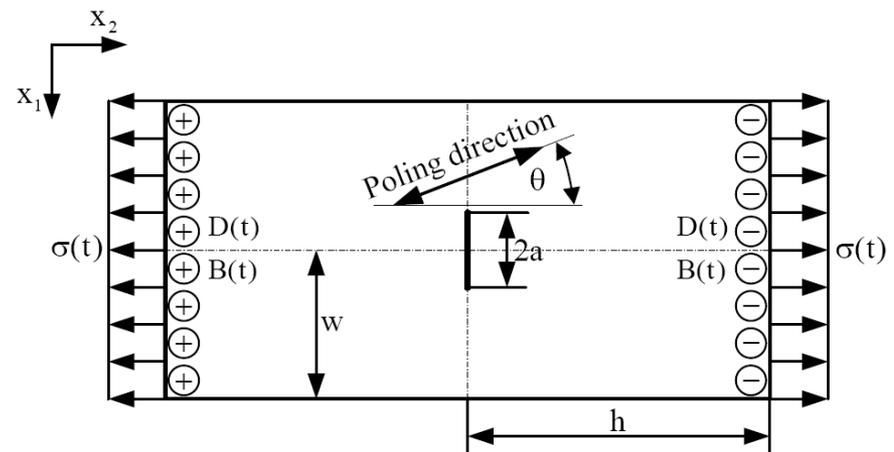


Some dynamic results: Finite cracked MEE plate under combined magneto-electro-mechanical impacts.

BaTiO₃
50%
CoFe₂O₄



$$\kappa = \frac{e_{22}}{\epsilon_{22}} \frac{D_0}{\sigma_0}, \quad \chi = \frac{h_{22}}{\gamma_{22}} \frac{B_0}{\sigma_0}$$

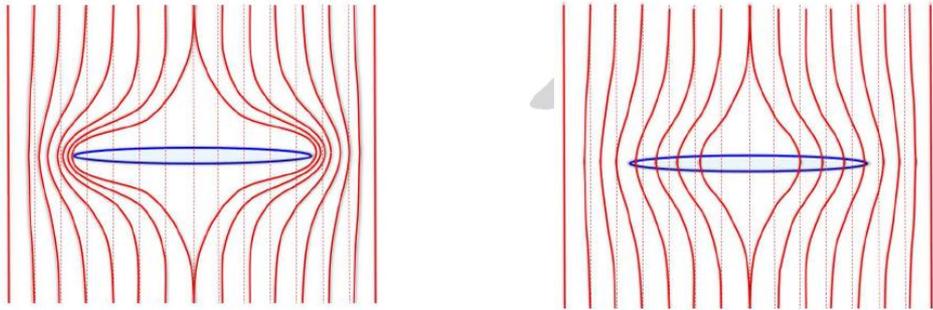


Concluding remarks

- ✓ BEM works for fracture applications!!, leading to accurate evaluation of the relevant fracture parameters

Some additional issues

- ❖ Realistic boundary conditions



- ❖ Include other relevant variables (T...) and material nonlinearities
- ❖ Develop adequate fracture criteria for multifield materials
- ❖ Improve fundamental solutions
- ❖ and much more...

Thanks for your attention!!

andres@us.es



Some references of our group's work:

Wünsche, M., García-Sánchez, F., Sáez, A., Zhang, Ch.
A 2D time-domain collocation-Galerkin BEM for dynamic crack analysis in piezoelectric solids
(2010) Engineering Analysis with Boundary Elements, 34 (4), pp. 377-387.

Rojas-Díaz, R., García-Sánchez, F., Sáez, A.
Analysis of cracked magnetoelastoelectric composites under time-harmonic loading
(2010) International Journal of Solids and Structures, 47 (1), pp. 71-80.

Rojas-Díaz, R., García-Sánchez, F., Sáez, A., Zhang, Ch.
Dynamic crack interactions in magnetoelastoelectric composite materials
(2009) International Journal of Fracture, 157 (1-2), pp. 119-130.

García-Sánchez, F., Zhang, C., Sáez, A.
2-D transient dynamic analysis of cracked piezoelectric solids by a time-domain BEM
(2008) Computer Methods in Applied Mechanics and Engineering, 197 (33-40), pp. 3108-3121.

Rojas-Díaz, R., Sáez, A., García-Sánchez, F., Zhang, C.
Time-harmonic Green's functions for anisotropic magnetoelastoelectricity
(2008) International Journal of Solids and Structures, 45 (1), pp. 144-158.

García-Sánchez, F., Rojas-Díaz, R., Sáez, A., Zhang, Ch.
Fracture of magnetoelastoelectric composite materials using boundary element method (BEM)
(2007) Theoretical and Applied Fracture Mechanics, 47 (3), pp. 192-204.

Sáez, A., García-Sánchez, F., Domínguez, J.
Hypersingular BEM for dynamic fracture in 2-D piezoelectric solids
(2006) Computer Methods in Applied Mechanics and Engineering, 196 (1-3), pp. 235-246.



Some references of our group's work:

García-Sánchez, F., Sáez, A., Domínguez, J.
Two-dimensional time-harmonic BEM for cracked anisotropic solids
(2006) Engineering Analysis with Boundary Elements, 30 (2), pp. 88-99.

García-Sánchez, F., Sáez, A., Domínguez, J.
Anisotropic and piezoelectric materials fracture analysis by BEM
(2005) Computers and Structures, 83 (10-11 SPEC. ISS.), pp. 804-820.

García, F., Sáez, A., Domínguez, J.
Traction boundary elements for cracks in anisotropic solids
(2004) Engineering Analysis with Boundary Elements, 28 (6), pp. 667-676.

Sáez, A., Domínguez, J.
Dynamic crack problems in three-dimensional transversely isotropic solids
(2001) Engineering Analysis with Boundary Elements, 25 (3), pp. 203-210.

Sáez, A., Domínguez, J.
Far field dynamic Green's functions for BEM in transversely isotropic solids
(2000) Wave Motion, 32 (2), pp. 113-123.

Sáez, A., Ariza, M.P., Domínguez, J.
Three-dimensional fracture analysis in transversely isotropic solids
(1997) Engineering Analysis with Boundary Elements, 20 (4), pp. 287-298.